

Computing the action ground states of nonlinear Schrödinger equation

Wei LIU

National University of Defense Technology, China

Collaborators: Chushan Wang, Tingfeng Wang, Zhenye Wen, Yongjun Yuan, Xiaofei Zhao

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Outline

- 1 Introduction
- 2 Computing action ground state in focusing case
- 3 Computing action ground state in defocusing case
- 4 Properties of defocusing action ground state
- 5 Conclusion

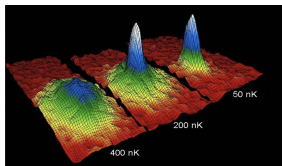
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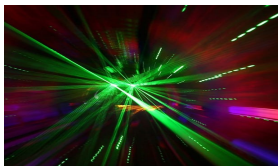
Rotating Nonlinear Schrödinger equation (RNLS)

$$i\partial_t\psi(\mathbf{x},t) = -\frac{1}{2}\Delta\psi + V(\mathbf{x})\psi + \beta|\psi|^{p-1}\psi - \Omega L_z\psi,$$

- $\psi : \mathbb{R}^d \times \mathbb{R}^+ \rightarrow \mathbb{C}$ unknown, $\mathbf{x} \in \mathbb{R}^d$, $t > 0$, $i = \sqrt{-1}$
- $V(\mathbf{x}) \geq 0$ external potential, e.g., $V(\mathbf{x}) = \frac{1}{2} \sum_{j=1}^d \gamma_j^2 x_j^2$
- $\beta > 0$ (defocusing) or $\beta < 0$ (focusing), $1 < p < \frac{d+2}{d-2}$ (if $d \geq 3$)
- $\Omega \in \mathbb{R}$ the rotational speed, $L_z = -i(x_1\partial_{x_2} - x_2\partial_{x_1})$ if $d \geq 2$; $L_z = 0$ if $d = 1$



Bose-Einstein condensate



Nonlinear optics



Deep water waves

Stationary Equation of RNLS

- Stationary state solutions (solitons):

$$\psi(\mathbf{x}, t) = e^{i\omega t} \phi(\mathbf{x}), \quad \omega \in \mathbb{R} \text{ (chemical potential),}$$

$$H(\phi) := -\frac{1}{2}\Delta\phi + V(\mathbf{x})\phi + \beta|\phi|^{p-1}\phi - \Omega L_z\phi + \omega\phi = 0, \quad \phi(\mathbf{x}) \not\equiv 0$$

- Nontrivial solutions of elliptic equation $-\Delta\phi = g(\phi)$ in $\mathbb{R}^d$:
 - Existence & Multiplicity: W. Strauss (1977); H. Berestycki & P.L. Lions (1983) ...
- Identify the relatively stable, physically relevant one [H. Berestycki & P.L. Lions (1983)]:
 - treat ω as a given parameter — action ground state
 - treat ω as a Lagrange multiplier — energy ground state

Definition of Action Ground State: H. Berestycki & P.L. Lions (1983)

- Prescribe a chemical potential $\omega \in \mathbb{R}$. Introduce the action functional:

$$S_{\Omega,\omega}(\phi) := \int_{\mathbb{R}^d} \left(\frac{1}{2} |\nabla \phi|^2 + V(\mathbf{x}) |\phi|^2 + \frac{2\beta}{p+1} |\phi|^{p+1} - \Omega \bar{\phi} L_z \phi + \omega |\phi|^2 \right) d\mathbf{x}$$

- Euler-Lagrange Equation: stationary states $\iff$ critical points of $S_{\Omega,\omega}(\phi)$,

$$\frac{\delta S_{\Omega,\omega}(\phi)}{\delta \bar{\phi}} = H(\phi) = \left(-\frac{1}{2} \Delta + V + \beta |\phi|^{p-1} - \Omega L_z + \omega \right) \phi = 0$$

Action Ground State ϕ_g (minimizes $S_{\Omega,\omega}$ among all the nontrivial solutions)

$$S_{\Omega,\omega}(\phi_g) = \min_{\phi \neq 0, H(\phi)=0} S_{\Omega,\omega}(\phi)$$

- $S_{\Omega,\omega}(\phi)$ is not bounded from below in H^1 in focusing case ($\beta < 0$)
- $H(\phi) = 0$ may have infinitely many solutions: H. Berestycki & P.L. Lions (1983) ...

Existing Related Works

- Existence of action ground state in **focusing case** ($\beta < 0$):
 - Non-rotating, no potential ($\Delta\phi + g(\phi) = 0$): Berestycki & Lions (1983) ...
 - Non-rotating ($\Omega = 0$) with harmonic potential $V(\mathbf{x}) = \frac{1}{2} \sum_{j=1}^d \gamma_j^2 x_j^2$: Fukuizumi (2001)
 - Rotating case ($\Omega \in \mathbb{R}$): Ardila & Hajaiej (2021)
 - ...

Existence & variational characterization in defocusing case ($\beta > 0$) is rarely discussed!

- Numerical methods for computing solitary waves (and multiple nontrivial solutions):
 - Fixed-point iteration type method: Petviashvili & Pokhotelov (1992); ...
 - minimax-type methods: Choi & McKenna (1993); Li & Zhou (2001); ...
 - saddle-point dynamics: E & Zhou (2011); Zhang & Du (2012); Yin, Zhang & Zhang (2019); ...
 - ...

No specific numerical techniques for computing action ground state ϕ_g of Rotating NLS so far.

Energy ground state

- Energy Ground State: For a given mass $m > 0$,

$$u_g \in \arg \min_{\|u\|_{L^2}^2 = m} E_\Omega(u),$$

where the energy functional [Bao, Wang & Markowich (2005); Bao & Cai (2013)]

$$E_\Omega(u) := \int_{\mathbb{R}^d} \left(\frac{1}{2} |\nabla u|^2 + V(\mathbf{x}) |u|^2 + \frac{2\beta}{p+1} |u|^{p+1} - \Omega \bar{u} L_z u \right) d\mathbf{x}$$

- Q: Connection between two kinds of ground states?

$$\begin{array}{ccccccc} \omega & \longrightarrow & \min_{\phi \neq 0, H(\phi)=0} & S_{\Omega, \omega}(\phi) & \longrightarrow & \phi_g & \longrightarrow & \|\phi_g\|_{L^2}^2 \\ \vdots & & & & & & & \vdots \\ \omega_g & \longleftarrow & u_g & \longleftarrow & \min_{\|u\|_{L^2}^2 = m} & E_\Omega(u) & \longleftarrow & m \end{array}$$

- Recent results in focusing non-rotating case: Dovetta et al. (2022); Jeanjean & Lu (2022)
- No results in defocusing case.

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Focusing case: Existence & variational characterization

- Nehari manifold $\mathcal{M} = \{\phi \in X \setminus \{0\} : K_{\Omega,\omega}(\phi) = 0\}$, where

$$K_{\Omega,\omega}(\phi) := \int_{\mathbb{R}^d} \bar{\phi} H(\phi) d\mathbf{x} = \int_{\mathbb{R}^d} \left(\frac{1}{2} |\nabla \phi|^2 + V |\phi|^2 - \Omega \bar{\phi} L_z \phi + \beta |\phi|^{p+1} + \omega |\phi|^2 \right) d\mathbf{x}$$

- Clearly, $\mathcal{M}$ contains all nontrivial solutions of $H(\phi) = 0$

Lemma 1

Let $\beta < 0$, $V(\mathbf{x}) = \frac{1}{2} \sum_{j=1}^d \gamma_j^2 x_j^2$, $\gamma_j \geq 0$. If $\omega > -\lambda_0$ (with λ_0 the smallest eigenvalue of $-\frac{1}{2}\Delta + V - \Omega L_z$) and $|\Omega| < \min\{\gamma_1, \gamma_2\}$, then there exists an action ground state ϕ_g and

$$S_{\Omega,\omega}(\phi_g) = \min_{\phi \in \mathcal{M}} S_{\Omega,\omega}(\phi).$$

cf. Berestycki & Lions (1983) ($V = 0, \Omega = 0$); Fukuizumi (2001) ($\Omega = 0$); Ardila & Hajaiej (2021)

Method I: Minimization with Nehari Constraint

- A continuous (projected) gradient flow to solve directly $\min_{\phi \in \mathcal{M}} S_{\Omega, \omega}(\phi)$:

$$\partial_t \phi(\mathbf{x}, t) = -\frac{\delta S_{\Omega, \omega}(\phi)}{\delta \phi} + \lambda \frac{\delta K_{\Omega, \omega}(\phi)}{\delta \phi} = -H(\phi) + \lambda \frac{\delta K_{\Omega, \omega}(\phi)}{\delta \phi}$$

where $\lambda(t)$ is chosen to preserve the constraint, i.e. $K_{\Omega, \omega}(\phi(\cdot, t)) = 0 \ \forall t > 0$.

- Action-diminishing property: $\frac{d}{dt} S_{\Omega, \omega}(\phi) = -2 \|\partial_t \phi\|_{L^2}^2$
- Efficient gradient flow schemes and preconditioned optimization methods can be designed [C. Wang (JCP, 2022)], [L-Wen-Yuan-Zhao (JCP 2025)]
- The formulation seems a bit complicated. Can we look for a simplification?

Method II: Minimization with L^{p+1} -normalization

Theorem 2 (L-Yuan-Zhao, SISC, 2023)

Let $\beta < 0$, $\omega > -\lambda_0$, and $|\Omega| < \min\{\gamma_1, \gamma_2\}$. Then

$$\phi_g \text{ is an action GS} \iff Q(u_*) = \min_{u \in \mathcal{S}_{p+1}} Q(u)$$

where $u_*(\mathbf{x}) := \phi_g(\mathbf{x}) / \|\phi_g\|_{L^{p+1}}$, $\mathcal{S}_{p+1} := \{u \in X : \|u\|_{L^{p+1}} = 1\}$ and

$$Q(u) := \int_{\mathbb{R}^d} \left(\frac{1}{2} |\nabla u|^2 + V(\mathbf{x})|u|^2 - \Omega \bar{u} L_z u + \omega |u|^2 \right) d\mathbf{x}$$

Moreover, $\phi_g(\mathbf{x}) = \left(\frac{Q(u_*)}{-\beta} \right)^{\frac{1}{p-1}} u_*(\mathbf{x})$.

A continuous normalized gradient flow

- A continuous normalized gradient flow for solving $\min_{u \in \mathcal{S}_{p+1}} Q(u)$:

$$\partial_t u(\mathbf{x}, t) = \left(\frac{1}{2} \Delta - V + \Omega L_z - \omega + \lambda(u(\cdot, t)) |u|^{p-1} \right) u, \quad u(\cdot, 0) = u_0 \in \mathcal{S}_{p+1},$$

$$\lambda(u) = \frac{\int_{\mathbb{R}^d} |u|^{p-1} \bar{u} \left(-\frac{1}{2} \Delta u + (V + \omega)u - \Omega L_z u \right) d\mathbf{x}}{\int_{\mathbb{R}^d} |u|^{2p} d\mathbf{x}}$$

- Normalization-conservation: $\frac{d}{dt} \|u\|_{L^{p+1}}^{p+1} = 0$
- Energy diminishing: $\frac{d}{dt} Q(u) = -2 \|\partial_t u\|_{L^2}^2$
- Drawbacks:
 - strong regularity requirement on u ;
 - hard to construct energy-decaying schemes

Gradient flow with asymptotic Lagrange multiplier

- Normalized gradient flow with asymptotic Lagrange multiplier (GFALM) [L-Yuan-Zhao (2023)]:

$$\begin{cases} \partial_t u = \left(\frac{1}{2} \Delta - V + \Omega L_z - \omega + \tilde{\lambda}(u(\cdot, t_n)) |u(\cdot, t_n)|^{p-1} \right) u, & t \in [t_n, t_{n+1}), \\ u(\mathbf{x}, t_{n+1}) := u(\mathbf{x}, t_{n+1}^+) = \frac{u(\mathbf{x}, t_{n+1}^-)}{\|u(\cdot, t_{n+1}^-)\|_{L^{p+1}}}, & n \geq 0, \quad u(\mathbf{x}, 0) = u_0(\mathbf{x}), \end{cases}$$

where $u_0 \in \mathcal{S}_{p+1}$ and $t_n = n\tau$ ($n \geq 0$) with $\tau > 0$ a time step.

- $\tilde{\lambda}$ is an **asymptotic Lagrange multiplier** defined as

$$\tilde{\lambda}(u) = \frac{Q(u)}{\int_{\mathbb{R}^d} |u|^{p+1} d\mathbf{x}}$$

In fact, $\tilde{\lambda}(u(\cdot, t_n)) = Q(u(\cdot, t_n))$ since $u(\cdot, t_n) \in \mathcal{S}_{p+1}$

- A similar framework to GFDN/GFLM for computing L^2 -normalized energy ground states:
Bao-Du (2004), Bao-Cai (2013), L-Cai (2021), ...

Backward-forward Euler scheme with stabilization

- **GFALM-BF:**

$$\begin{cases} \frac{u^{(1)} - u^n}{\tau} = \left(\frac{1}{2} \Delta - \alpha_n \right) u^{(1)} + \left(\alpha_n - V + \Omega L_z - \omega + \tilde{\lambda}^n |u^n|^{p-1} \right) u^n, \\ u^{n+1} = \frac{u^{(1)}}{\|u^{(1)}\|_{L^{p+1}}}, \quad n \geq 0, \quad u^0 = u_0 \in \mathcal{S}_{p+1}, \end{cases}$$

where $\tilde{\lambda}^n = \tilde{\lambda}(u^n) = Q(u^n)$ and $\alpha_n \geq 0$ is a stabilization factor.

- **Fixed point:** for $\forall \tau > 0, \alpha_n \geq 0$,

$$u^{n+1} = u^n \iff (\tilde{\lambda}^n, u^n) \text{ solves Euler-Lagrange eqn, i.e.,} \\ \left(-\frac{1}{2} \Delta + V - \Omega L_z + \omega \right) u^n = \tilde{\lambda}^n |u^n|^{p-1} u^n$$

- Only needs to solve a linear elliptic equation with **constant** coefficients
- In practice, we **truncate $\mathbb{R}^d$ to a bounded (box) domain $U \subset \mathbb{R}^d$** and impose the homogeneous Dirichlet BCs or periodic BCs (FFT solver).

Unconditional energy-decaying property

In the following, we consider the GFALM-BF scheme on a suitably large box domain $U \subset \mathbb{R}^d$ with the homogeneous Dirichlet or periodic BCs.

Theorem 3 (L-Yuan-Zhao, SISC 2023; L-Wang-Zhao, arXiv:2602.20820)

Assume that $V \in L^\infty(U)$ and

$$\alpha_n \geq \frac{1}{2} \max \left\{ 0, V(\mathbf{x}) + \omega + \frac{|\Omega|^2}{2} (x_1^2 + x_2^2) \right\} + \frac{1}{2}.$$

Then, GFALM-BF scheme has the *unconditionally energy-decaying property*: for $\forall \tau > 0$,

$$Q(u^{n+1}) - Q(u^n) \leq - \left(\frac{1}{2} \|u^{n+1} - u^n\|_{H^1}^2 + \frac{2}{\tau} \|u^{n+1} - u^n\|_{L^2}^2 \right) / \|u^{(1)}\|_{L^{p+1}}^2.$$

Theorem 4 (L-Wang-Zhao, arXiv:2602.20820)

Suppose that the above Assumptions hold. Let $u^0 \in H^1(U) \cap \mathcal{S}_{p+1}$. Then, for every fixed $\tau > 0$:

- 1 the solution $\{u^n\}_{n=0}^\infty$ of GFALM-BF has a strongly convergent subsequence $\{u^{n_j}\}_{j=0}^\infty$ in $H^1(U)$;
- 2 every accumulation point $u^* \in \mathcal{S}_{p+1}$ of $\{u^n\}_{n=0}^\infty$ in $H^1(U)$ satisfies the Euler-Lagrange equation:

$$\langle \mathcal{A}u^*, v \rangle = \lambda^* \langle |u^*|^{p-1}u^*, v \rangle, \quad \forall v \in H^1(U),$$

with the Lagrange multiplier $\lambda^* := \tilde{\lambda}(u^*) = Q(u^*) = \inf_{n \geq 0} Q(u^n)$. Here $\mathcal{A}u = \frac{\delta Q}{\delta u}$.

Local exponential convergence rate

Theorem 5 (L-Wang-Zhao, arXiv:2602.20820)

Suppose that the above basic Assumptions hold, and let u_g be a **non-degenerate** action GS.

Then, there exists a $\delta_0 > 0$ such that, for any initial guess $u^0 \in U_{\delta_0}(u_g) \cap \mathcal{S}_{p+1}$ and any $\tau > 0$, the GFALM-BF solution $\{u^n\}$ converges to u_g at the exponential rate:

$$\|u^n - u_g\|_{H^1} \leq C e^{-an\tau/(1+\tau)}, \quad \forall n \in \mathbb{N},$$

where $C, a > 0$ are constants independent of n and τ .

Convergence to discrete GS

- Fourier pseudospectral discretization:

$$u_{g,h} \in \operatorname{argmin}_{\|u_h\|_{\ell^{p+1}}=1} Q_h(u_h), \quad Q_h(u_h) := \frac{1}{2} \langle (-D_{xx}^F) u_h, u_h \rangle_h + h \sum_{j=0}^{M-1} (V_j + \omega) |u_j|^2.$$

Theorem 6 (L-Wang-Zhao, arXiv:2602.20820)

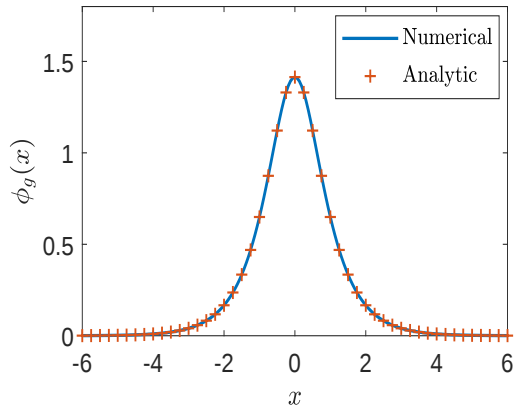
For any fixed h , assume that the discrete GS $u_{g,h}$ satisfies the **non-degeneracy** condition.

Then, there exists a $\delta_0 > 0$ such that, for any initial data $u_h^0 \in U_{\delta_0}(u_{g,h}) \cap \mathcal{S}_{h,p+1}$ and any $\tau > 0$, the fully discrete GFALM scheme converges at the rate:

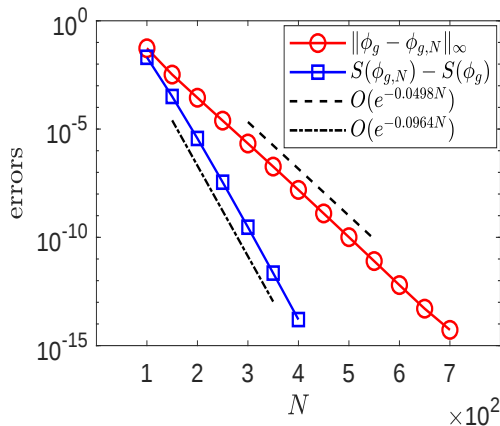
$$\|u_h^n - u_{g,h}\|_{1,h} \leq C e^{-an\tau/(1+\tau)}, \quad \forall n \in \mathbb{N}.$$

Spatial spectral accuracy of GFALM-BF

$$\frac{1}{2}\partial_{xx}\phi(x) + |\phi(x)|^2\phi(x) = \phi(x), \quad \phi_g(x) = \sqrt{2}\operatorname{sech}(\sqrt{2}x), \quad x \in \mathbb{R}$$



(a) Profiles of the numerical solution $\phi_{g,N}$ with $N = 2^{10}$ and the analytic solution ϕ_g



(b) Error of the numerical solution with respect to N

Exponential convergence rate

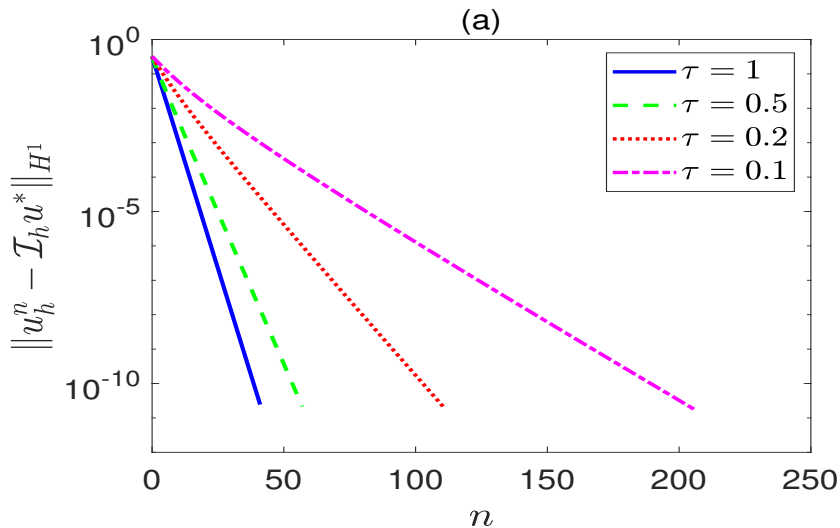


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Existence & variational characterization in defocusing case

Theorem 7 (L-Yuan-Zhao, SISC 2023; L-Wen-Yuan-Zhao, JCP 2025)

Let $\beta > 0$, $\omega < -\lambda_0$. Assume that $V(\mathbf{x}) \geq 0$ ($\forall \mathbf{x} \in \mathbb{R}^d$), $\lim_{|\mathbf{x}| \rightarrow \infty} V(\mathbf{x}) = \infty$, and

$$|\Omega| < \gamma_* := \sup \left\{ \gamma > 0 : \lim_{|\mathbf{x}| \rightarrow \infty} \left(V(\mathbf{x}) - \frac{\gamma^2}{2} (x_1^2 + x_2^2) \right) = \infty \right\} \text{ (if } d \geq 2 \text{)}.$$

Then, $S_{\Omega, \omega}$ is bounded from below in X and there exists $\phi_g \in X$ such that

$$S_{\Omega, \omega}(\phi_g) = \inf_{\phi \in X} S_{\Omega, \omega}(\phi) = \inf_{\phi \in \mathcal{M}} S_{\Omega, \omega}(\phi).$$

Moreover, ϕ_g is an action GS $\iff u_* := \phi_g / \|\phi_g\|_{L^{p+1}} \in \arg \min_{u \in \mathcal{S}_{p+1}} Q(u)$.

- Various numerical algorithms can be designed based on each of **three minimization formulations** [L-Wen-Yuan-Zhao (JCP 2025)]

Unconstrained minimization methods

- A direct gradient flow for unconstrained minimization:

$$\partial_t \phi(\mathbf{x}, t) = -\frac{\delta S_{\Omega, \omega}(\phi)}{\delta \bar{\phi}} = \frac{1}{2} \Delta \phi - V \phi - \beta |\phi|^{p-1} \phi + \Omega L_z \phi - \omega \phi, \quad t \geq 0,$$

- **Backward-forward Euler scheme (GF-BF):**

$$\frac{\phi^{n+1} - \phi^n}{\tau} = \frac{1}{2} \Delta \phi^{n+1} - \alpha_n \phi^{n+1} + (\alpha_n - V - \omega - \beta |\phi^n|^{p-1} + \Omega L_z) \phi^n,$$

- Unconditionally action-decaying property (for suitably large α_n)
- Convergence analysis in rotating case: [L-Wang-Yuan-Zhao (arXiv:2605.04485)]
- Preconditioned optimization methods ... [L-Yuan-Zhao (SISC/2023)]; [L-Wen-Yuan-Zhao (JCP/2025)]

Unconditionally decay of action

Theorem 8 (L-Wang-Yuan-Zhao, arXiv:2605.04485)

Let U be a box domain, $V \in L^\infty(U)$, $\beta > 0$ and $\phi^0 \in X$. There exists a constant $C_0 = C_0(U, V, \beta, p, \phi_0)$ s.t., when $\alpha \geq \frac{1}{2} \max \left\{ 0, \operatorname{ess\,sup}_{\mathbf{x} \in U} (V(\mathbf{x}) + \omega + 2|\Omega|^2(x_1^2 + x_2^2)) \right\} + C_0$, the GF-BF scheme

$$\frac{\phi^{n+1} - \phi^n}{\tau} = \frac{1}{2} \Delta \phi^{n+1} - \alpha \phi^{n+1} + (\alpha - V - \omega - \beta |\phi^n|^{p-1} + \Omega L_z) \phi^n,$$

possesses the unconditional decaying property for the action functional $S_{\Omega, \omega}$: $\forall \tau > 0$,

$$S_{\Omega, \omega}(\phi^{n+1}) - S_{\Omega, \omega}(\phi^n) \leq -\frac{2}{\tau} \|\phi^{n+1} - \phi^n\|_{L^2}^2 - \frac{1}{4} \|\phi^{n+1} - \phi^n\|_{H^1}^2.$$

Based on this property, we can further establish a global convergence and a local exponential convergence rate to the action ground state under a reasonable non-degeneracy condition [L-Wang-Yuan-Zhao (arXiv:2605.04485)].

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Main attention: Relation with energy ground states

- **Energy Ground State** $u_g = u_g(\mathbf{x})$ minimizes the energy [Bao & Cai (2013)]

$$E_{\Omega}(u) := \int_{\mathbb{R}^d} \left(\frac{1}{2} |\nabla u|^2 + V(\mathbf{x}) |u|^2 + \frac{2\beta}{p+1} |u|^{p+1} - \Omega \bar{u} L_z u \right) d\mathbf{x}$$

under the constraint $\|u\|_{L^2}^2 = m$ with **given mass** $m > 0$.

- The chemical potential ω_g serves as a Lagrange multiplier:

$$\omega_g = -\mu(u_g), \quad \mu(u) := \frac{1}{m} \int_{\mathbb{R}^d} \left(\frac{1}{2} |\nabla u|^2 + V(\mathbf{x}) |u|^2 + \beta |u|^{p+1} - \Omega \bar{u} L_z u \right) d\mathbf{x}$$

- **Q: Connection between two kinds of ground states?**

$$\begin{array}{ccccccc} \omega & \longrightarrow & \min_{\phi \neq 0, H(\phi)=0} S_{\Omega, \omega}(\phi) & \longrightarrow & \phi_g & \longrightarrow & \|\phi_g\|_{L^2}^2 \\ \vdots & & & & & & \vdots \\ \omega_g & \longleftarrow & u_g & \longleftarrow & \min_{\|u\|_{L^2}^2=m} E_{\Omega}(u) & \longleftarrow & m \end{array}$$

Equivalence in non-rotating case

In the following, we focus on the defocusing case ($\beta > 0$).

Theorem 9 (L-Wang-Zhao, M3AS, 2025)

Let $\Omega = 0$. Then *two kinds of ground states are equivalent*:

- For any $\omega < -\lambda_0$, an action ground state ϕ_ω associated with ω is also an energy ground state with mass $m = \|\phi_\omega\|_{L^2}^2$.
- For any $m > 0$, an energy ground state u_g with mass m is also an action ground state associated with $\omega = -\mu(u_g)$ where $\mu(u) := \frac{1}{\|u\|_{L^2}^2} \int_{\mathbb{R}^d} \left(\frac{1}{2} |\nabla u|^2 + V(\mathbf{x})|u|^2 + \beta|u|^{p+1} \right) d\mathbf{x}$.

Conditional equivalence in rotating case

- In rotating case ($\Omega \neq 0$), due to the **non-uniqueness**, we consider

$$\phi_{\Omega,\omega} \in \mathcal{A}_{\Omega,\omega} := \arg \min_{\phi \in X} S_{\Omega,\omega}(\phi), \quad \phi_m^E \in \mathcal{B}_{\Omega,m} := \arg \min_{\|\phi\|_{L^2}^2 = m} E(\phi),$$

Theorem 10 (L-Wang-Zhao, M3AS, 2025)

For any $\omega < -\lambda_0(\Omega)$, the *action ground state* $\phi_{\Omega,\omega}$ is also an *energy ground state* with mass $m = \|\phi_{\Omega,\omega}\|_{L^2}^2$, and *all the energy ground states* ϕ_m^E with precisely this mass m are also *action ground states* associated with ω . Moreover,

$$\mathcal{A}_{\Omega,\omega} = \bigcup_{m \in \mathcal{N}_{\Omega,\omega}} \mathcal{B}_{\Omega,m}, \quad \text{where } \mathcal{N}_{\Omega,\omega} := \{\|\phi\|_{L^2}^2 : \phi \in \mathcal{A}_{\Omega,\omega}\}.$$

- **Q: Is there $m > 0$ that is not the mass of any action ground state?**

Numerical exploration: Vortices and critical angular velocity

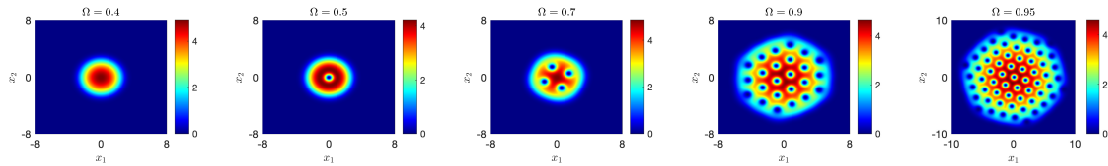


Figure 1: Density profiles of $|\phi_{\Omega, \omega}(x_1, x_2)|^2$ for $\omega = -5$ and different Ω .

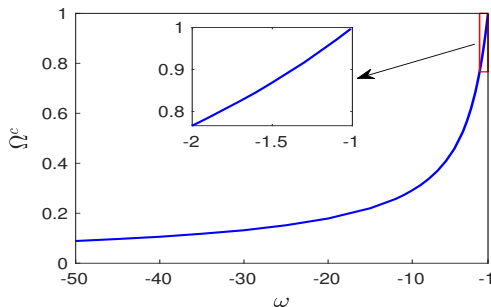
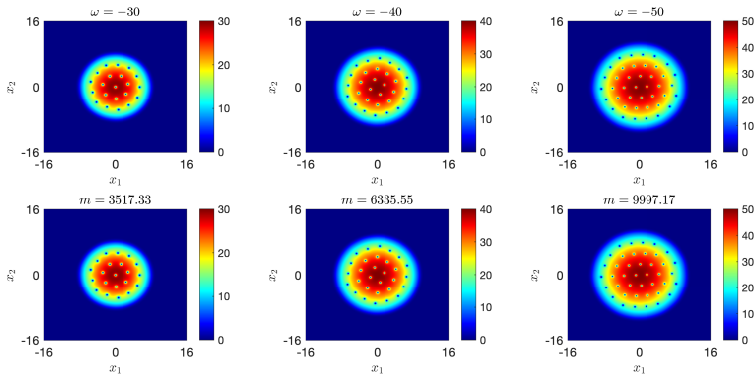


Figure 2: Critical angular velocity Ω^c as a function of ω

Examples of equivalence in rotating case ($\Omega = 0.5$)

Action ground states (1st row) and corresponding energy ground states (2nd row)



ω	S_g	m	E_g	μ_g	e_ω^{rel}	e_S^{rel}
-30	-34342.56	3517.33	71177.33	30.00	3.21×10^{-10}	7.31×10^{-14}
-40	-82943.09	6335.55	170478.80	40.00	6.59×10^{-10}	4.82×10^{-14}
-50	-163950.62	9997.17	335907.93	50.00	2.21×10^{-10}	2.38×10^{-14}

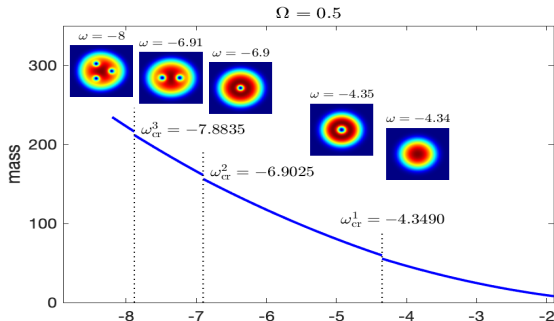
Numerical evidence of non-equivalence: Jump of mass

- Observation: the mass of the action ground state $\|\phi_{\Omega,\omega}\|_{L^2}^2$ has **jump discontinuities at** $\omega_{\text{cr}}^1 = -4.3490$, $\omega_{\text{cr}}^2 = -6.9025$, $\omega_{\text{cr}}^3 = -7.8835$, etc.

- No action ground state with a mass value coming from the range

$$(55.4535, 59.6198) \cup (156.3416, 161.2327) \cup (211.8255, 216.2082) \cup \dots$$

- ω_{cr}^j are exactly the transition points where the number of vortices in $\phi_{\Omega,\omega}$.



Numerical evidence: Energy GSs that are not action GSs

- For $m \in (55.4535, 59.6198)$, the energy ground states ϕ_m^E (1st row) are not any action ground state (2nd row), and there exists a critical value $m_{\text{cr}}^1 \in (57.494, 57.495)$ s.t.
 - when $m \in (55.4535, m_{\text{cr}}^1)$, $\omega = -\mu(\phi_m^E) < \omega_{\text{cr}}^1$ and $\|\phi_{\Omega, \omega}\|_{L^2}^2 > 59.6198$;
 - when $m \in (m_{\text{cr}}^1, 59.6198)$, $\omega = -\mu(\phi_m^E) > \omega_{\text{cr}}^1$ and $\|\phi_{\Omega, \omega}\|_{L^2}^2 < 59.6198$.
- These energy ground states are nontrivial solutions to the stationary RNLS with $\omega = -\mu(\phi_m^E)$ and **higher action values**, e.g. for $m = 57.494$, $S_{\Omega, \omega}(\phi_m^E) = -77.4181 > S_g = -77.7370$

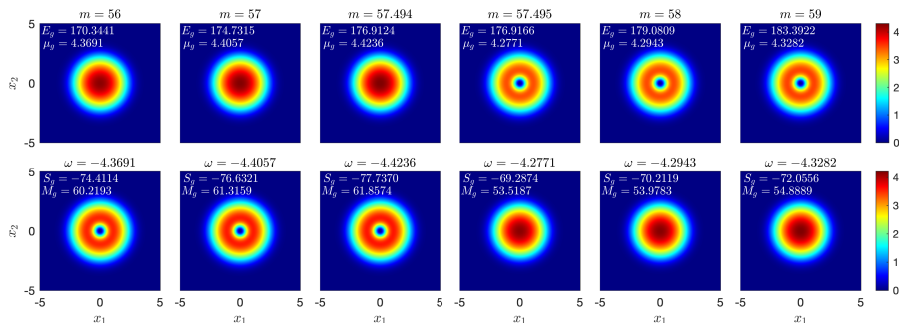


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Conclusion

In focusing case ($\beta < 0$):

- $\phi_g \iff$ Nehari minimization $\iff$ minimization with L^{p+1} -normalization
- Gradient flow with asymptotic Lagrange multiplier (GFALM) and its convergence

In defocusing case ($\beta > 0$):

- Existence, variational characterization, minimization algorithms:
 $\phi_g \iff$ Nehari minimization $\iff$ L^{p+1} -constrained min $\iff$ unconstrained min.
- Action GS vs. Energy GS:
 - Rotating case: Conditional equivalence & Numerical evidence for non-equivalence
 - Non-rotating case: complete equivalence

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-  W. Liu, C. Wang, X. Zhao. [On action ground states of defocusing nonlinear Schrödinger equations](#). *Math. Models Methods Appl. Sci.*, **35** (2025), pp. 39-74.
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-  W. Liu, T. Wang, and X. Zhao. [Convergence analysis of \$L^{p+1}\$ -normalized gradient flow for action ground state of nonlinear Schrödinger equation](#), arXiv:2602.20820 [math.NA]
-  W. Liu, T. Wang, Y. Yuan, and X. Zhao. [Analysis of gradient flow for computing defocusing action ground states of rotating nonlinear Schrödinger equations](#), arXiv:2605.04485 [math.NA]

Thank you!